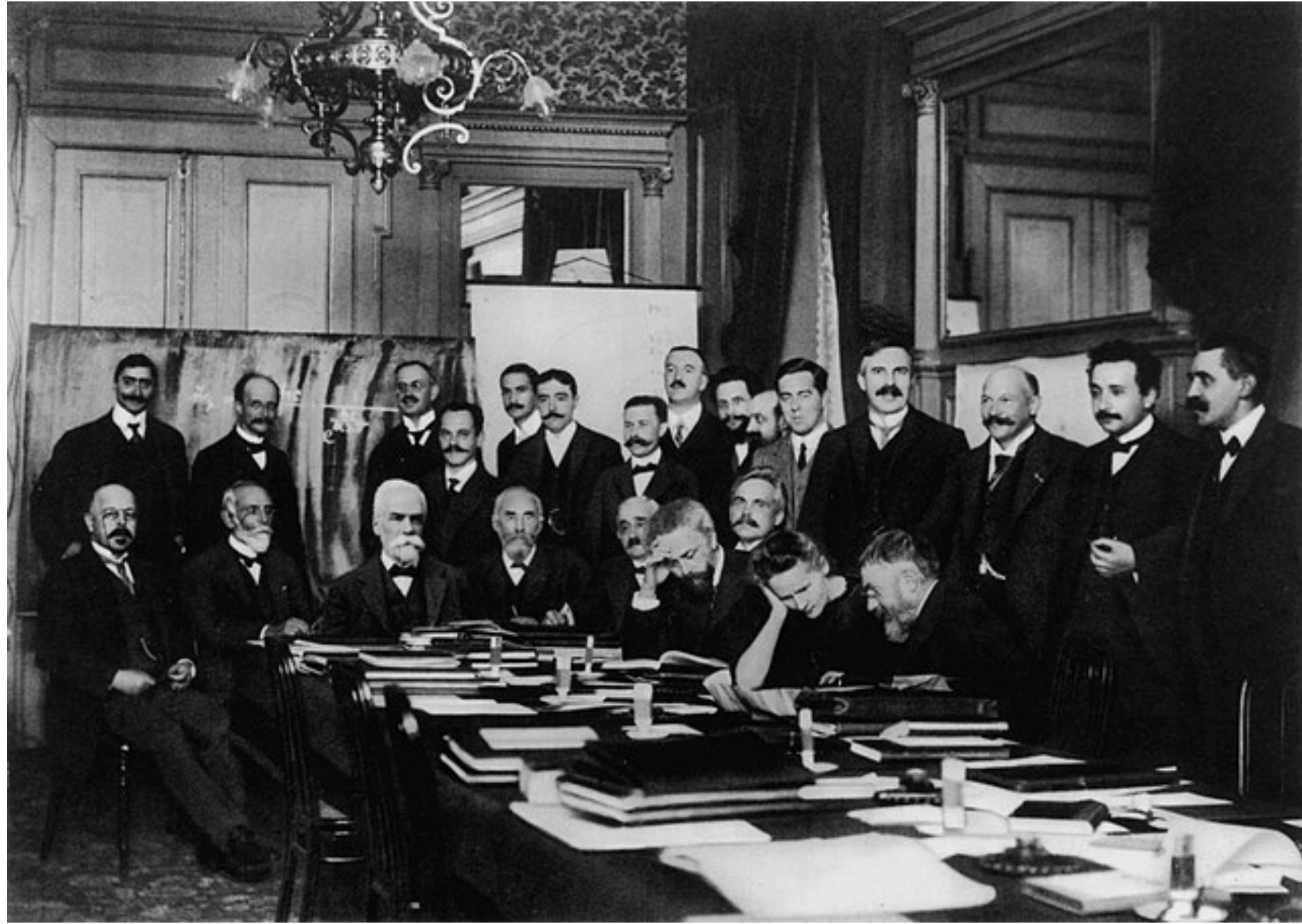


# **Non-linear Quantum Mechanics**

# Non Linear Quantum Mechanics?



**Theory built on observations in the 1900s  
Why should it be “the absolute truth”?**

## What?

**Two Postulates of Quantum Mechanics**

**Probability**

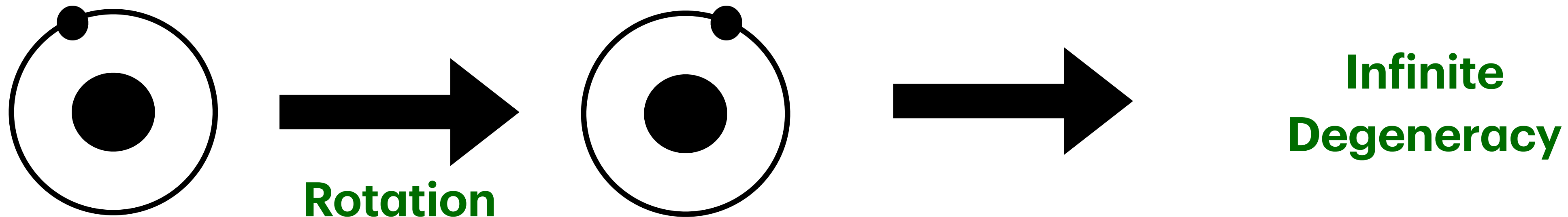
**Linearity**

## Which?

# Probability

Finite system has a finite set of energies  
Continuous observables and symmetries } Deterministic Observables?

Could an electron in an atom have a well defined position?



Quantum Mechanics

Sacrifice Determinism.

Preserve finite set of energy states, continuous symmetries and observables

Bell Inequalities, Kochen-Specker, SSC Theorems

# Causality and Entanglement

**Trial Non-Linear Term**

$$i \frac{\partial \Psi}{\partial t} = H_L \Psi + \epsilon (\Psi^2 + \Psi^{*2}) \Psi$$

**Entanglement is fundamental to quantum mechanics**

$$\Psi(x, y; t) = \sum_{i,j} c_{ij}(t) \alpha_i(x) \beta_j(y)$$

**Apply some local operation on x:  $a_i(x) \rightarrow U a_i(x)$**

**Does it instantly change the time evolution of y?**

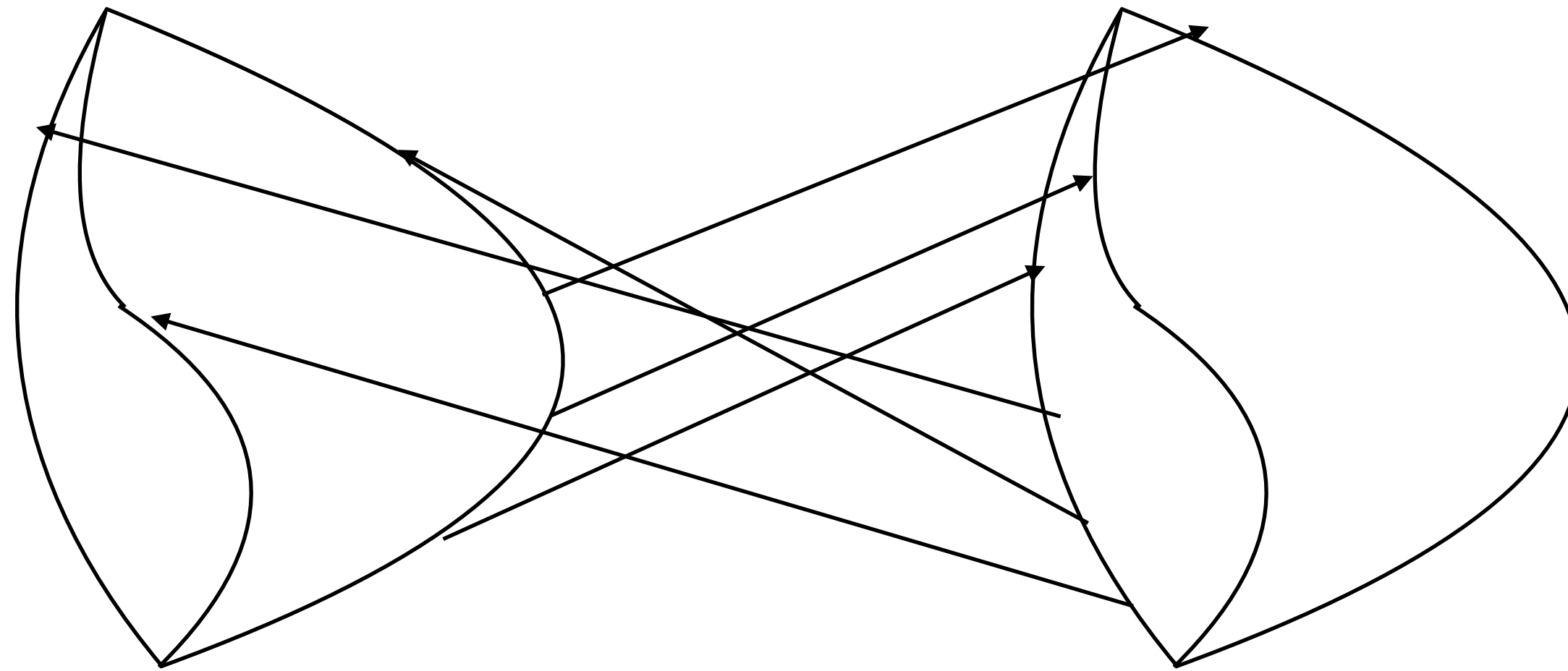
**YES**

**Not causal**

# Causality and Non-Linearity

## Linear Quantum Mechanics

### Electron Coupled to Electromagnetism



**Electron paths do not  
interact via  
electromagnetism**

**Paths of two electrons  
interact causally (QFT)**

**Why can't path talk to itself?  
Formulate directly into QFT**

# The Framework

## The Schrodinger Picture of Quantum Field Theory

$|\chi(t)\rangle$

**Quantum State of Fields**  
(e.g. in Fock states)

$\phi(x)$

**Time Independent**  
**Operators**

$$H = \int d^3x \mathcal{H}(\phi(x), \pi(x))$$

### Time Evolution

$$i \frac{\partial |\chi(t)\rangle}{\partial t} = H |\chi(t)\rangle$$

### Action

$$S = \int dt (i \langle \chi | \dot{\chi} \rangle - \langle \chi | H | \chi \rangle)$$



# The Framework

**Yukawa**  $H \supset \int d^3x \, y \, \phi(x) \, \bar{\Psi}(x) \, \Psi(x)$

## Action

$$S = \int dt \, (i \langle \chi | \dot{\chi} \rangle - \langle \chi | H | \chi \rangle) \supset \langle \chi(t) | \left( \int d^3x \, y \, \phi(x) \, \bar{\Psi}(x) \, \Psi(x) \right) | \chi(t) \rangle$$
$$\supset \left( \int d^3x \, y \, \langle \chi(t) | \phi(x) \, \bar{\Psi}(x) \, \Psi(x) | \chi(t) \rangle \right)$$

**Quantum Field Theory**  $\supset \left( \int d^3x \, y \, \langle \chi(t) | \phi(x) \, \bar{\Psi}(x) \, \Psi(x) + \frac{\phi^2}{\Lambda} \bar{\Psi} \Psi + \dots | \chi(t) \rangle \right)$

**Non-linearities in the operators but not in the state**

# The Framework

**Yukawa**  $H \supset \int d^3x \, y \, \phi(x) \, \bar{\Psi}(x) \, \Psi(x)$

**Linear QFT:**  $S \supset \left( \int d^3x \, y \, \langle \chi(t) | \phi(x) \, \bar{\Psi}(x) \, \Psi(x) | \chi(t) \rangle \right)$

**Non-Linear QFT:**  $S_{NL} \supset \epsilon \left( \int d^3x \, \langle \chi(t) | \phi(x) | \chi(t) \rangle \langle \chi(t) | \bar{\Psi}(x) \, \Psi(x) | \chi(t) \rangle \right)$

**Obeys all the rules**

**Higher order in states - leads to state dependent quantum evolution**

**Analyze non-linearity perturbatively**



# Perturbation Theory

$$\mathcal{H} \supset y\Phi\bar{\Psi}\Psi = (y\phi + \epsilon\langle\chi|\phi|\chi\rangle)\bar{\Psi}\Psi$$

$$i\frac{\partial|\chi\rangle}{\partial t} = H|\chi\rangle$$

**At zeroth order, this is just standard QFT**

**At first order, use zeroth order solution - expectation value is simply a background field**

**Perform standard QFT on this background field to compute first order correction**

**Applies to all orders : To compute term of given order, only need lower order terms**

**Lower order terms enter as background fields**

**Causality: Non-linearity enters via expectation value. At lowest order, causal from QFT.**

**Causal background field for all higher orders**

# Gauge Theories and Gravitation

Linear QFT Lagrangian, Shift bosonic field by expectation value

To Path Integral, add:

$$e^{iS_0 + i \int d^4x \left( e^{\frac{(A_\mu + \epsilon_\gamma \langle \chi | A_\mu | \chi \rangle)}{1 + \epsilon_\gamma}} J^\mu + \epsilon_{\tilde{\gamma}} \langle \chi | F_{\mu\nu} | \chi \rangle F^{\mu\nu} \right)}$$

Background Field

Gravitation

$$e^{iS_0 + i \int d^4x (\epsilon_G \langle \chi | g_{\mu\nu} | \chi \rangle \partial^\mu \phi \partial^\nu \phi)}$$

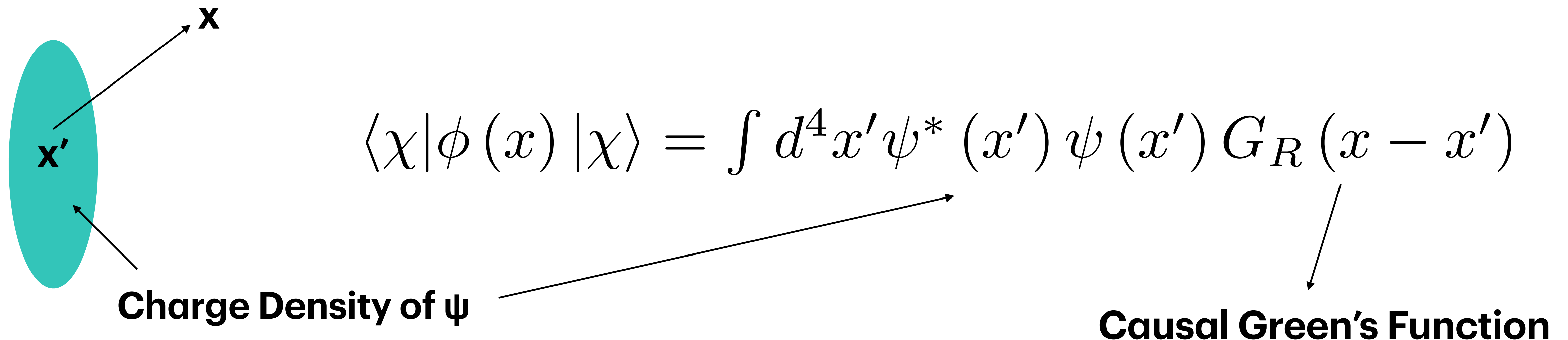
# Single Particle

$$\mathcal{L} \supset y \Phi \bar{\Psi} \Psi = y (\phi + \tilde{\epsilon} \langle \chi | \phi | \chi \rangle) \bar{\Psi} \Psi$$

Suppose we have a  $\psi$  particle - how does its wave-function evolve?

To zeroth order,  $\psi$  just sources the  $\Phi$  field

**Straightforward Computation of Expectation Value**


$$\langle \chi | \phi(x) | \chi \rangle = \int d^4 x' \psi^*(x') \psi(x') G_R(x - x')$$

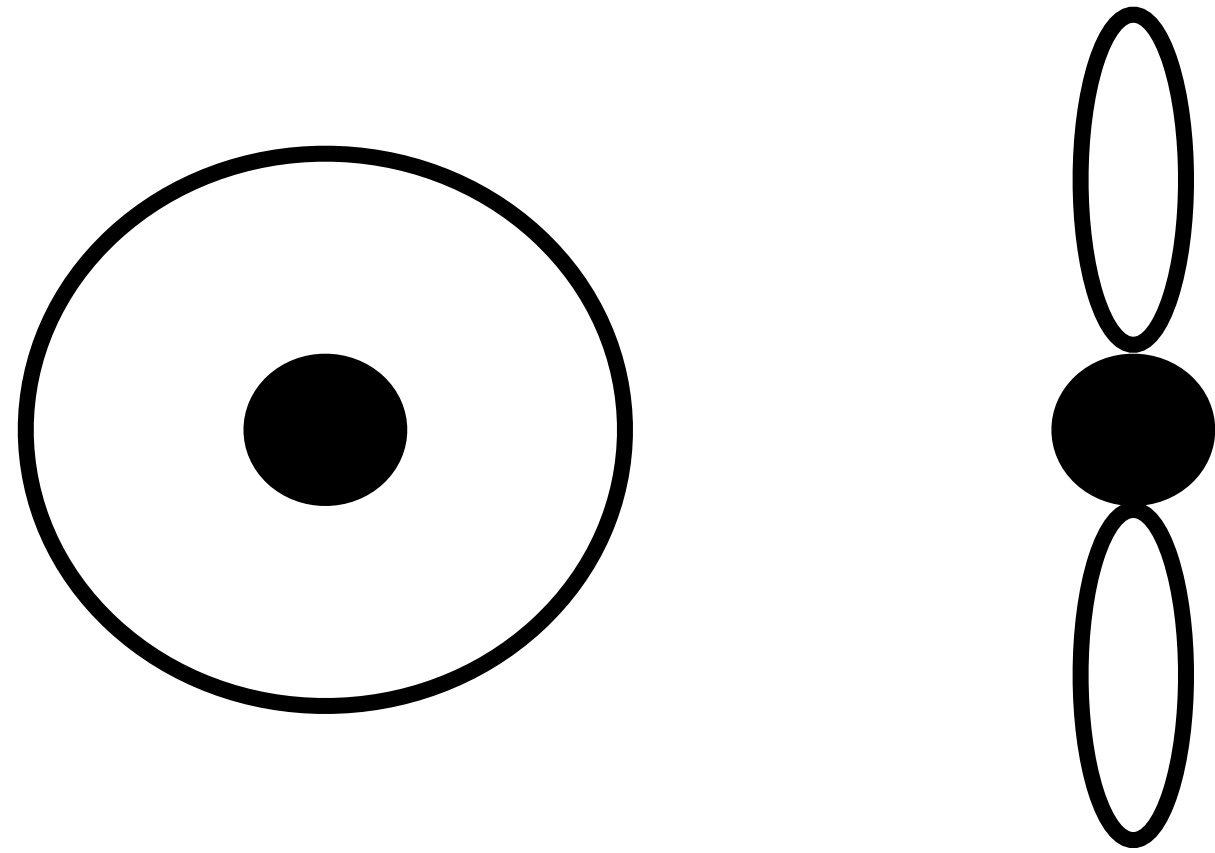
**Charge Density of  $\psi$**

**Causal Green's Function**

# Constraints

# Constraints

## What does this do to the Lamb Shift?



Proton at Fixed Location

2S and 2P electron have different charge distribution

Different expectation value of electromagnetic field

Level Splitting!

$$\langle \chi | A_\mu | \chi \rangle J^\mu$$

**BUT: Cannot decouple center of mass and relative co-ordinates**

Proton wave-function spread over some region (e.g. trap size ~ 100 nm)

Expectation value of electromagnetic field diluted

In neutral atom - heavily suppressed, except at edges!

$$\varepsilon < 10^{-2}$$

Similarly, kills possible bounds on QCD and gravity

# Experimental Tests



# Experimental Tests

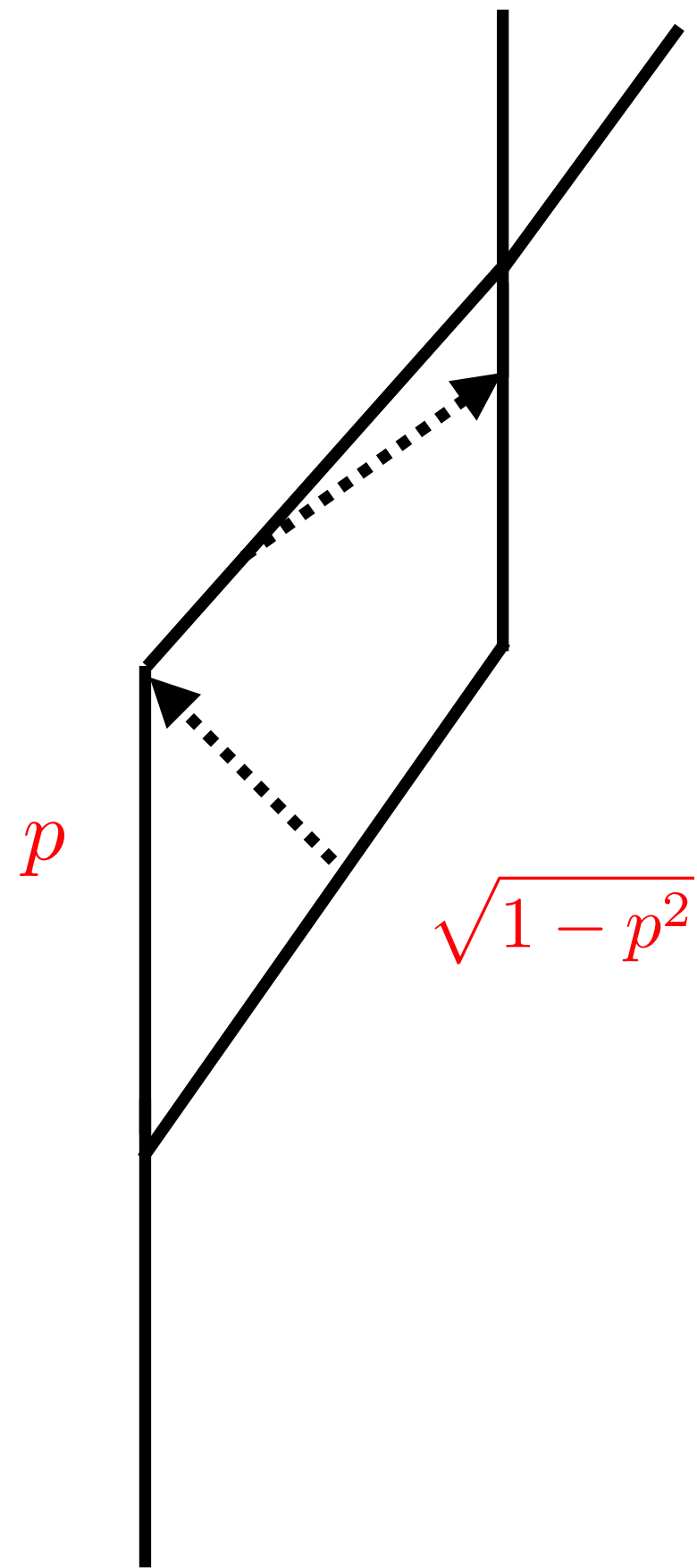
**Interferometry - interaction between paths**

**Take an ion - split its wave-function**

**Coulomb Field of one path interacts with the other path**

**Gives rise to phase shift that depends on the intensity  $p$  of the split**

**Use intensity dependence to combat systematics**

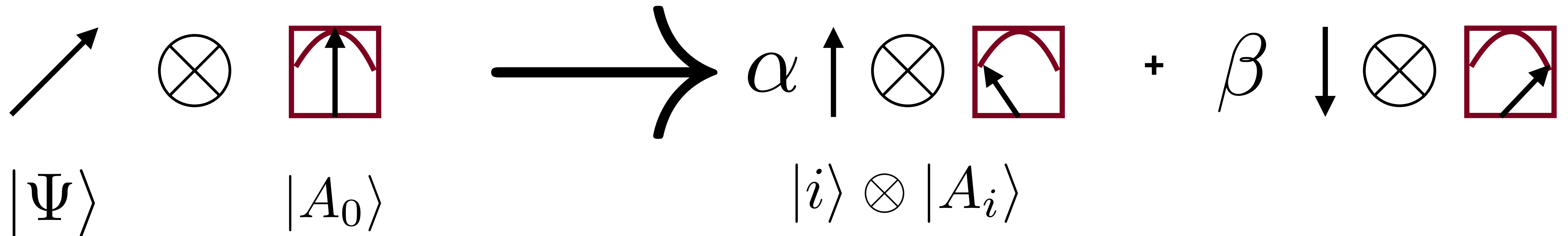


# Macroscopic Effects

# Measurement in Quantum Mechanics

Not some mysterious process

Interaction between quantum state and measuring device



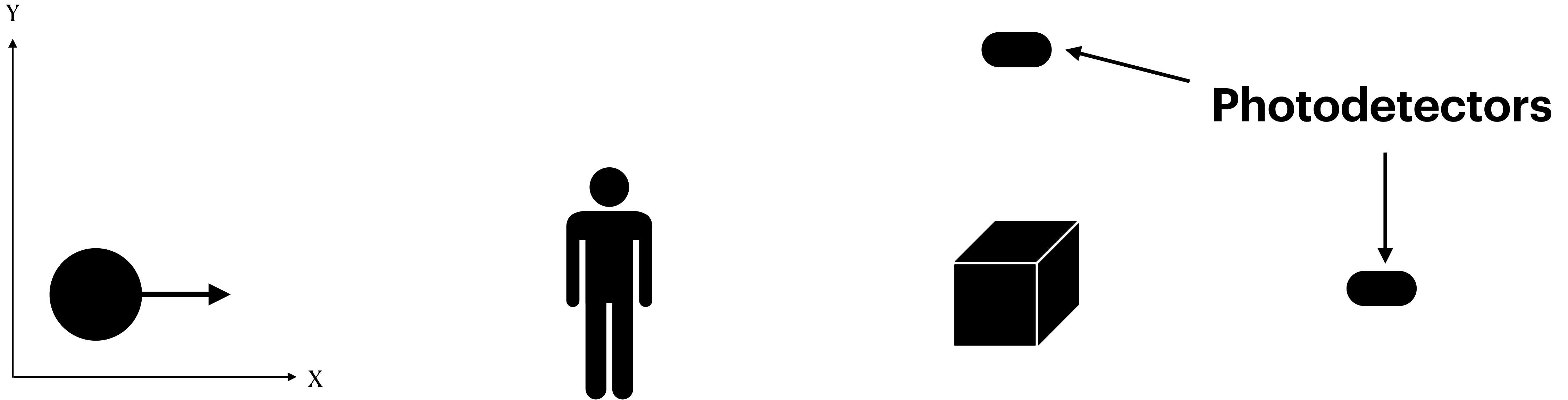
$$|\Psi\rangle \otimes |A_0\rangle \rightarrow \sum_i c_i |i\rangle \otimes |A_i\rangle$$

Prediction of Quantum Mechanics ("Many Worlds"), Not an interpretation

Pick:  $\langle A_j | A_i \rangle = \delta_{ij} \implies \rho_{|\Psi\rangle} = \sum_i c_i c_i^* |i\rangle \langle i|$

"Interpret" as direct  
sum of "worlds"

# Linear Quantum Mechanics



**Spin  
Along x**

**Experimentalist**

**Laser**

**Initial State :  $|\chi(0)\rangle$**

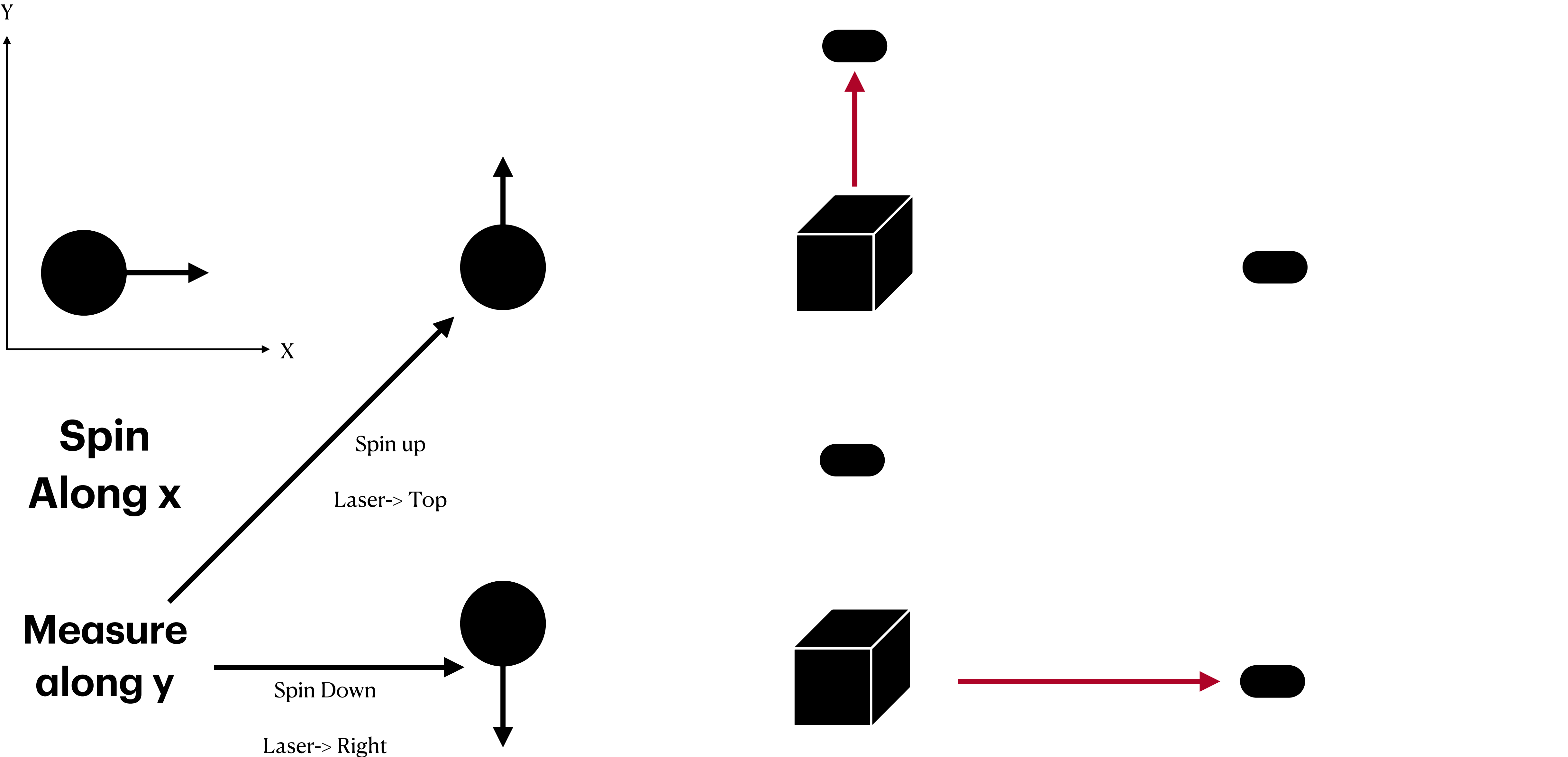
**Represents Full Quantum State (spin, experimentalist...)**

**Goal: Create Macroscopic Superposition**

**Method: Measure spin along y.**

**Depending upon outcome, send laser along different directions**

# Macroscopic Superposition



**Final State:**  $|X\rangle = |U\rangle|T\rangle|E_T\rangle + |D\rangle|R\rangle|E_R\rangle$

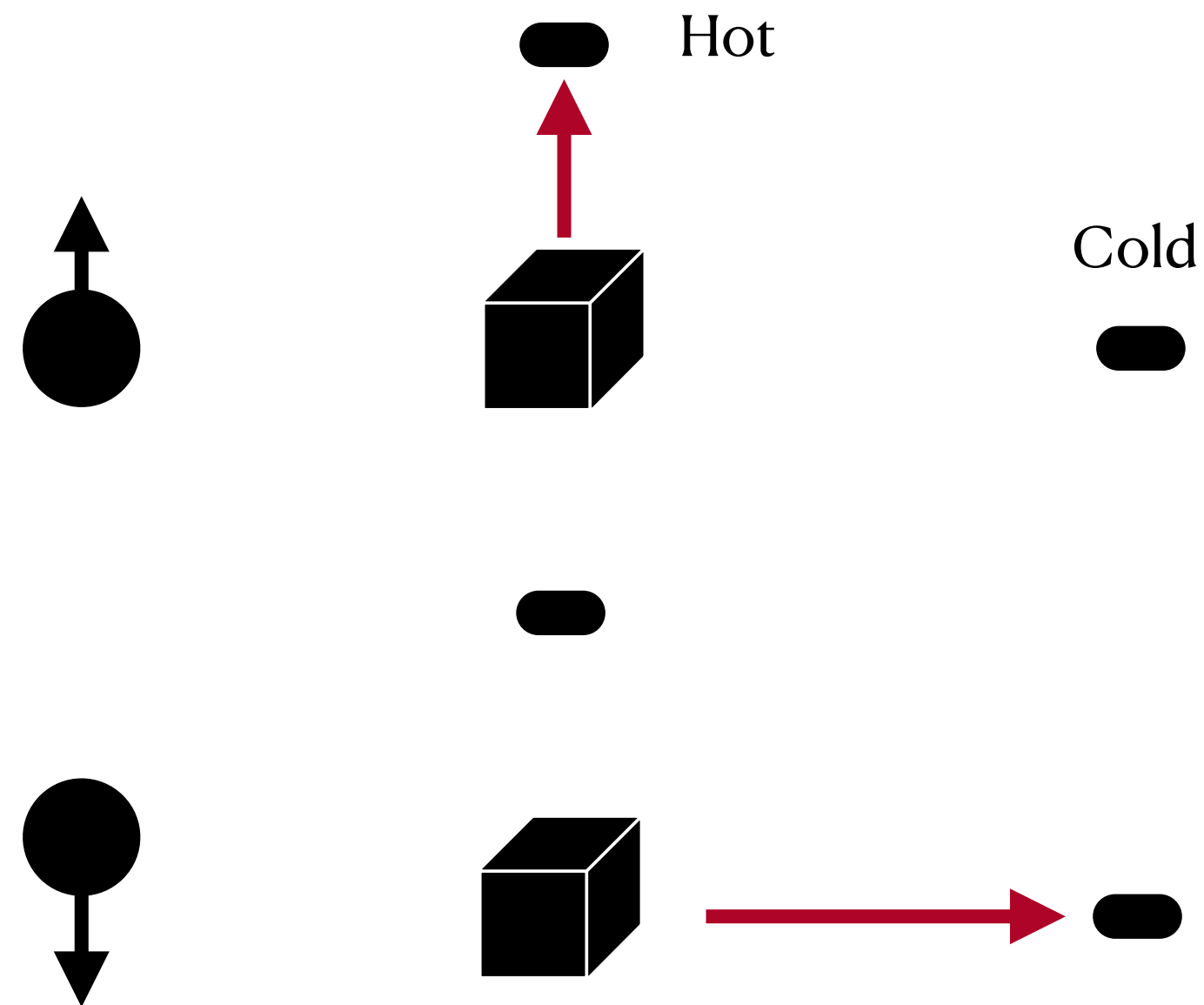
**Prediction of QM**

# Linear Quantum Mechanics

Which photodetectors light up?

$$|\chi\rangle = |U\rangle|T\rangle|E_T\rangle + |D\rangle|R\rangle|E_R\rangle$$

$$\mathcal{L} \supset eA_\mu \bar{\Psi} \gamma^\mu \Psi$$



Compute Transition Matrix Elements

$$\langle U | \langle T | \langle E_T | eA_\mu (x_T) \bar{\Psi} (x_T) \gamma^\mu \Psi (x_T) | \chi \rangle \neq 0$$

$$\langle U | \langle T | \langle E_T | eA_\mu (x_R) \bar{\Psi} (x_R) \gamma^\mu \Psi (x_R) | \chi \rangle = 0$$



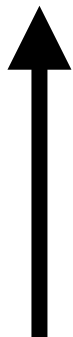
$$\langle T | A_\mu (x_R) | T \rangle = 0$$

But in both  $|E_T\rangle, |E_R\rangle$ :  $\langle \chi | A_\mu (x_T) | \chi \rangle \neq 0, \langle \chi | A_\mu (x_R) | \chi \rangle \neq 0$




# Non-Linear Quantum Mechanics


$$\mathcal{L} \supset e A_\mu \bar{\Psi} \gamma^\mu \Psi + \epsilon_\gamma e \langle \chi | A_\mu | \chi \rangle \bar{\Psi} \gamma^\mu \Psi$$



State Dependent Non-linear Term

But in both  $|E_T\rangle, |E_R\rangle$ :

  $\epsilon_{\text{Hot}}$   $\langle \chi | A_\mu (x_T) | \chi \rangle \neq 0, \langle \chi | A_\mu (x_R) | \chi \rangle \neq 0$

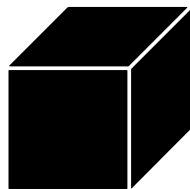
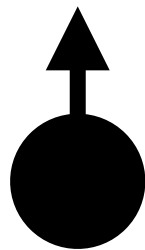
  $\epsilon_{\text{Hot}}$

Communication between “worlds”


Consequence of Causality - trace over entangled particles

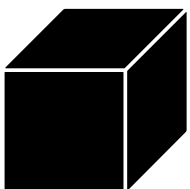
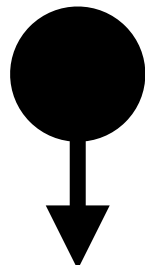
Non-linearity visible despite Environmental De-coherence!

Polchinski: “Everett Phone”



 Hot

  $\epsilon_{\text{Hot}}$



Hot



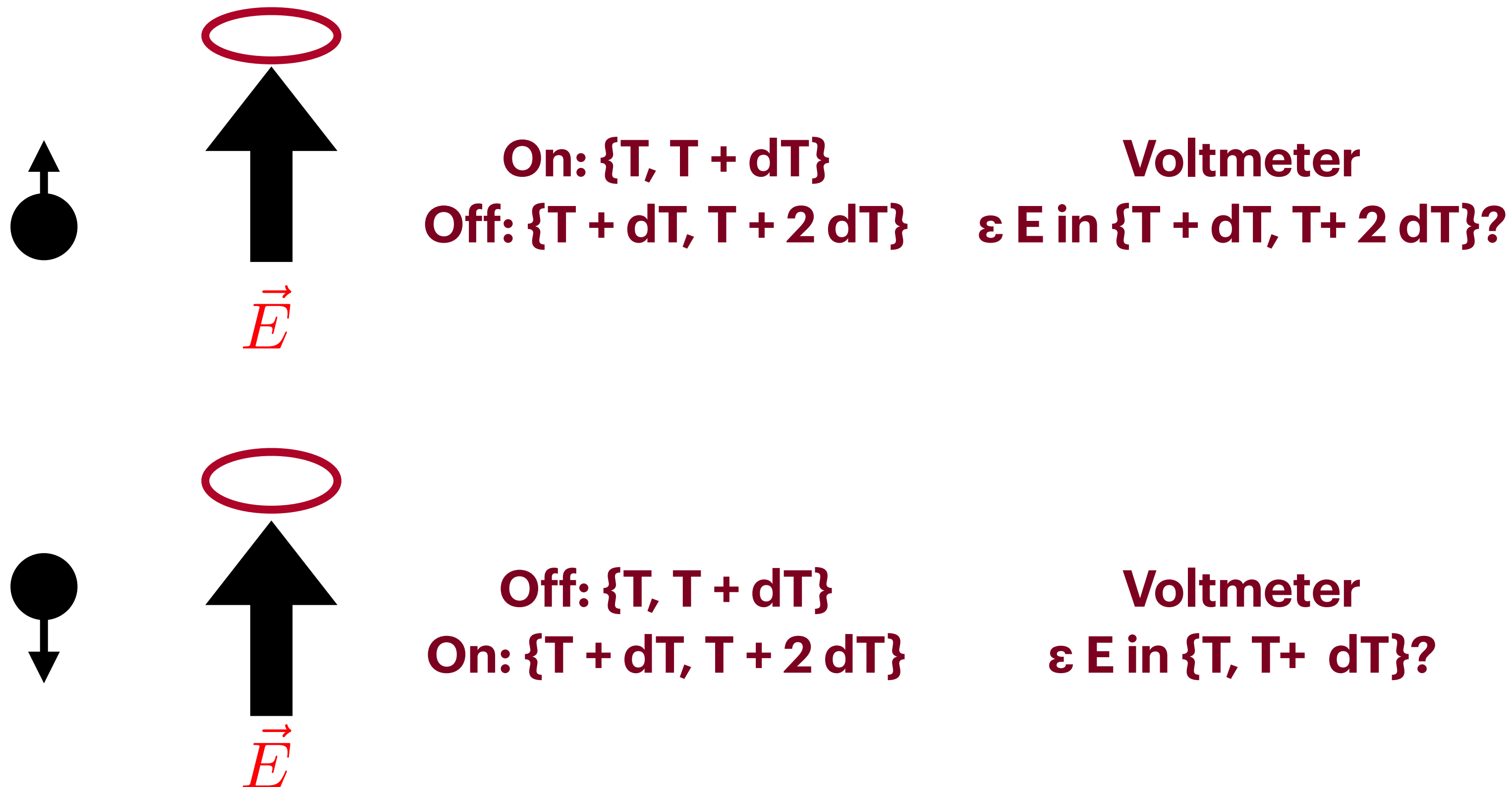
# Tests

# Experimental Tests

**Key Point: Create macroscopic superposition**

**Create Expectation value of EM/Gravity**

**Search for Expectation value**

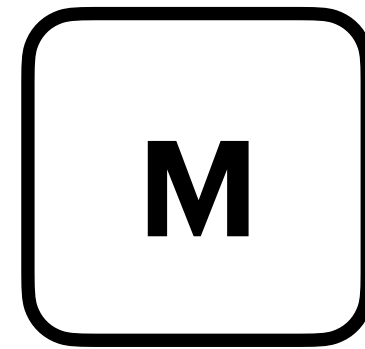
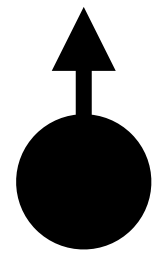


# Experimental Tests

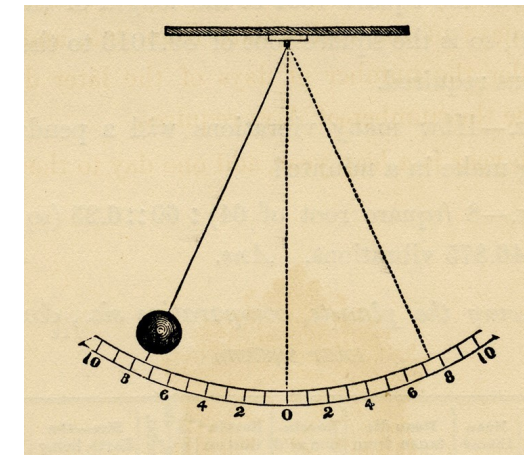
**Key Point: Create macroscopic superposition**

**Create Expectation value of EM/Gravity**

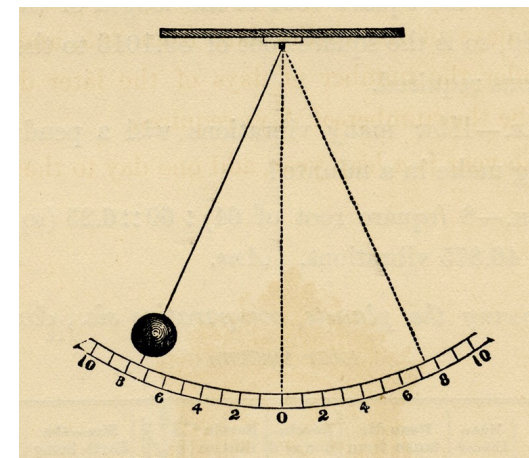
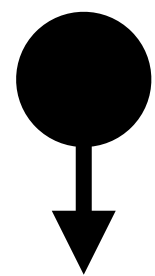
**Search for Expectation value**



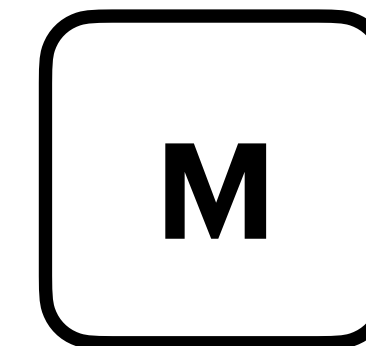
**X<sub>1</sub>**



**X<sub>2</sub>**



**X<sub>1</sub>**



**X<sub>2</sub>**

**Even Null Result is Interesting:**

$$G_{\mu\nu} = \langle T_{\mu\nu} \rangle$$

# Conclusions

- 1. Quantum Field Theory can be generalized to include non-linear, state dependent time evolution**
- 2. Conventional tests of quantum mechanics in atomic and nuclear systems do NOT probe causal non-linear quantum mechanics**
- 3. Straightforward set of experimental tests possible to probe non-linear quantum mechanics**
- 4. Motivation to test other extensions as well - e.g. Lindblad Decoherence**

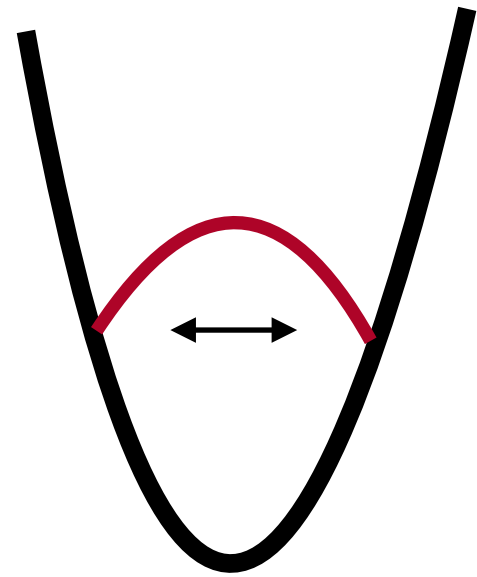
**Backup**



# Constraints

## Leading Constraint?

For  $\varepsilon > 0$  (repulsive interaction)

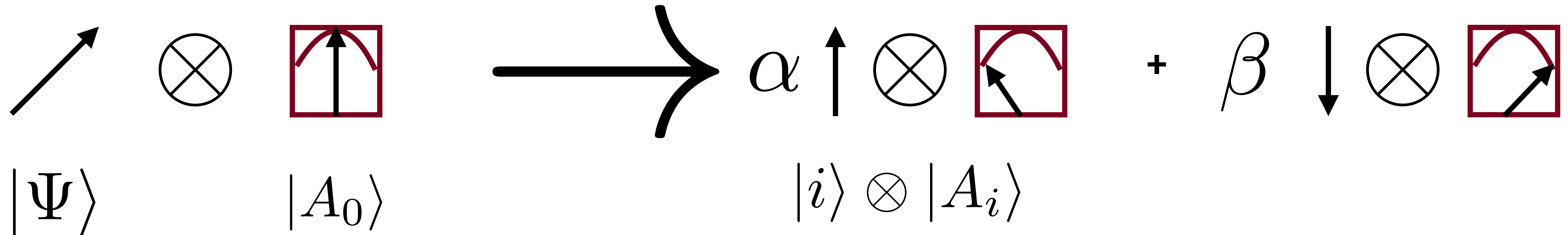


Too large a repulsion, Cant trap ion in trap  
 $\varepsilon < 10^{-5}$

No direct limit on  $\varepsilon < 0$  (attractive interaction)  
Perhaps from mapping of ion in trap?

# Measurement in Non-Linear Quantum Mechanics

Interaction between quantum state and measuring device



In linear QM, just need to know the basis vectors

Interaction Hamiltonian independent of unknown quantum state

$$\text{Pick: } \langle A_j | A_i \rangle = \delta_{ij}$$

**Key Point: Non-linear Hamiltonian depends upon unknown quantum state**

$$\text{No Guarantee: } \langle A_j | A_i \rangle = 0$$

$$|\Psi\rangle \otimes |A_0\rangle \rightarrow \sum_i c_i |i\rangle \otimes |A_i\rangle + \epsilon \sum_{i,j} d_{i,j} |i\rangle \otimes |A_j\rangle$$

**Measurement Noise**

# Framework

# The Framework

## The Schrodinger Picture of Quantum Field Theory

$|\chi(t)\rangle$

**Quantum State of Fields**  
(e.g. in Fock states)

$\phi(x)$

**Time Independent**  
**Operators**

$$H = \int d^3x \mathcal{H}(\phi(x), \pi(x))$$

### Time Evolution

$$i \frac{\partial |\chi(t)\rangle}{\partial t} = H |\chi(t)\rangle$$

### Action

$$S = \int dt i (\langle \chi | \dot{\chi} \rangle - \langle \chi | H | \chi \rangle)$$

# The Framework

**Yukawa**  $H \supset \int d^3x \, y \, \phi(x) \, \bar{\Psi}(x) \, \Psi(x)$

## Action

$$S = \int dt \, i \left( \langle \chi | \dot{\chi} \rangle - \langle \chi | H | \chi \rangle \right) \supset \langle \chi(t) | \left( \int d^3x \, y \, \phi(x) \, \bar{\Psi}(x) \, \Psi(x) \right) | \chi(t) \rangle$$
$$\supset \left( \int d^3x \, y \, \langle \chi(t) | \phi(x) \, \bar{\Psi}(x) \, \Psi(x) | \chi(t) \rangle \right)$$

**Quantum Field Theory**  $\supset \left( \int d^3x \, y \, \langle \chi(t) | \phi(x) \, \bar{\Psi}(x) \, \Psi(x) + \frac{\phi^2}{\Lambda} \bar{\Psi} \Psi + \dots | \chi(t) \rangle \right)$

**Non-linearities in the operators but not in the state**

# The Framework

**Yukawa**  $H \supset \int d^3x \, y \, \phi(x) \, \bar{\Psi}(x) \, \Psi(x)$

**Linear QFT:**  $S \supset \left( \int d^3x \, y \, \langle \chi(t) | \phi(x) \bar{\Psi}(x) \Psi(x) | \chi(t) \rangle \right)$

**Non-Linear QFT:**  $S_{NL} \supset \epsilon \left( \int d^3x \, \langle \chi(t) | \phi(x) | \chi(t) \rangle \langle \chi(t) | \bar{\Psi}(x) \Psi(x) | \chi(t) \rangle \right)$

**Obeys all the rules**

**Higher order in states - leads to state dependent quantum evolution**

**Analyze non-linearity perturbatively**



# Perturbation Theory

$$\mathcal{H} \supset y\Phi\bar{\Psi}\Psi = (y\phi + \epsilon\langle\chi|\phi|\chi\rangle)\bar{\Psi}\Psi$$

$$i\frac{\partial|\chi\rangle}{\partial t} = H|\chi\rangle$$

**At zeroth order, this is just standard QFT**

**At first order, use zeroth order solution - expectation value is simply a background field**

**Perform standard QFT on this background field to compute first order correction**

**Applies to all orders : To compute term of given order, only need lower order terms**

**Lower order terms enter as background fields**

**Causality: Non-linearity enters via expectation value. At lowest order, causal from QFT.**

**Causal background field for all higher orders**

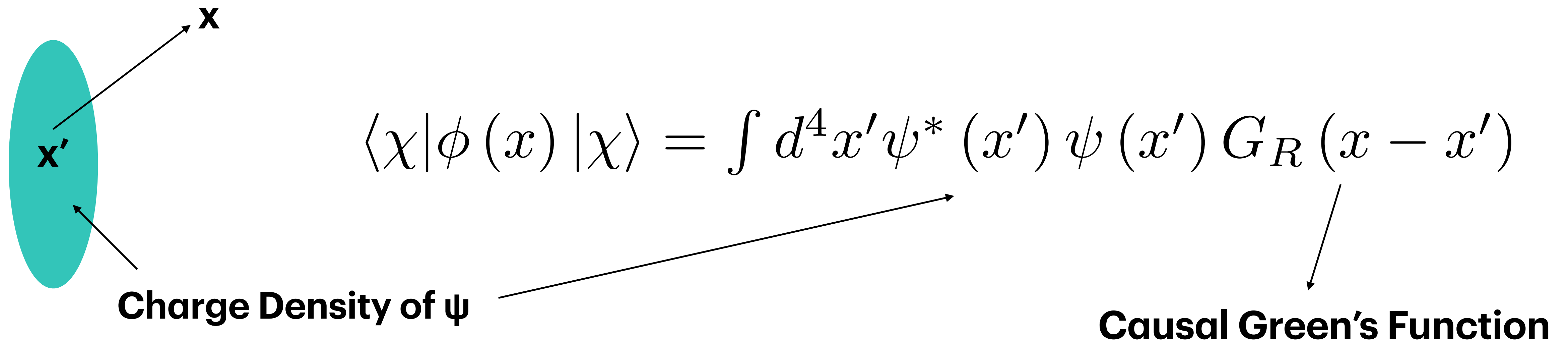
# Single Particle

$$\mathcal{L} \supset y \Phi \bar{\Psi} \Psi = y (\phi + \tilde{\epsilon} \langle \chi | \phi | \chi \rangle) \bar{\Psi} \Psi$$

Suppose we have a  $\psi$  particle - how does its wave-function evolve?

To zeroth order,  $\psi$  just sources the  $\Phi$  field

**Straightforward Computation of Expectation Value**


$$\langle \chi | \phi(x) | \chi \rangle = \int d^4 x' \psi^*(x') \psi(x') G_R(x - x')$$

**Charge Density of  $\psi$**

**Causal Green's Function**

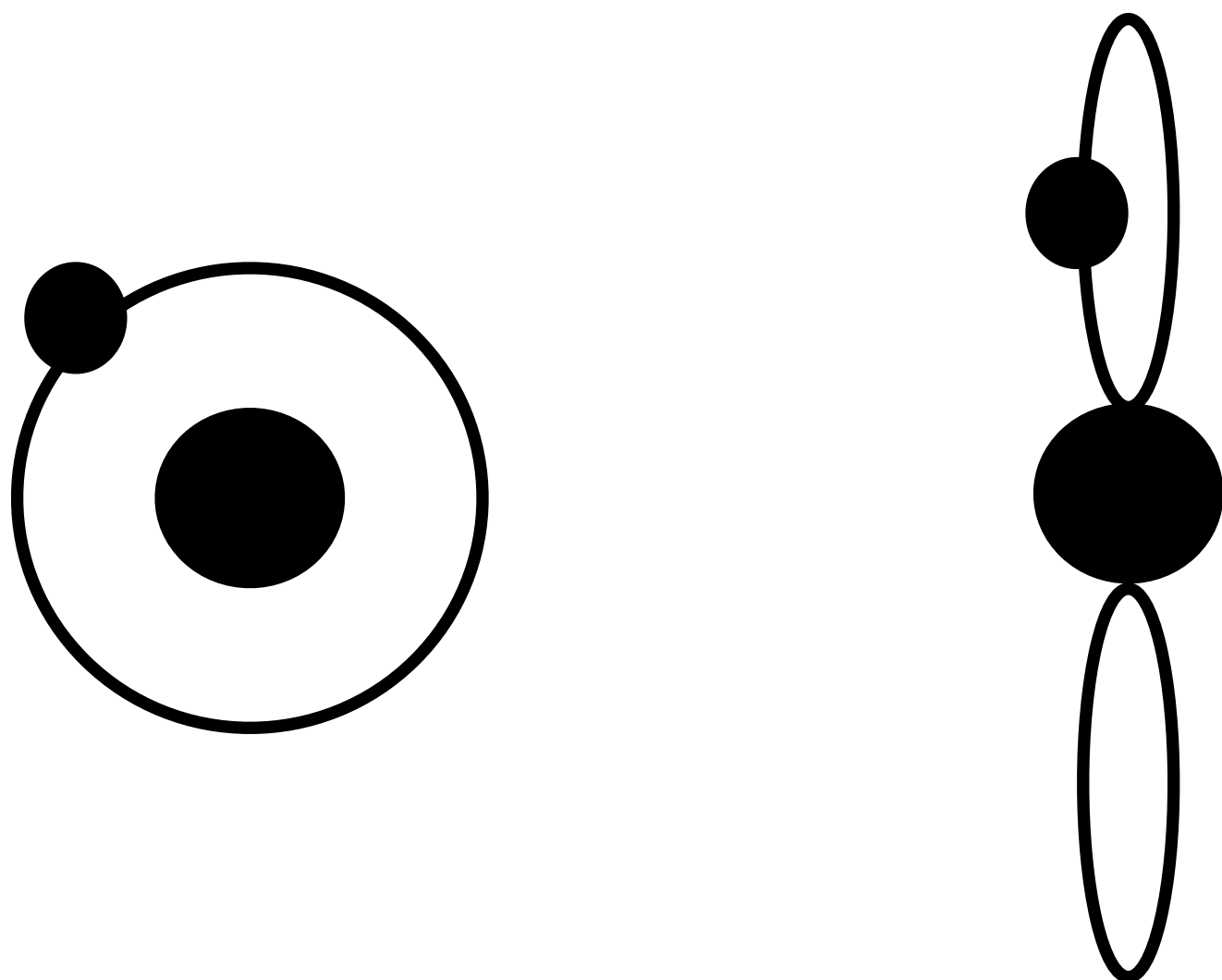
# Schrodinger Equation

$$\mathcal{H} \supset y\Phi\bar{\Psi}\Psi = (y\phi + \epsilon\langle\chi|\phi|\chi\rangle)\bar{\Psi}\Psi$$

**Single particle equation derived from field theory**

**Equation depends upon theory (Yukawa,  $\Phi^4$  etc)**

$$i\frac{\partial\Psi(t,\mathbf{x})}{\partial t} = \left(H + \tilde{\epsilon}y \int d^4x' \Psi^*(x) \Psi(x') G_R(x; x')\right) \Psi(t, \mathbf{x})$$



**Fixed Central particle**

**Self interaction of wave-function breaks degeneracy of levels**

**Hermitean Form of Hamiltonian implies conserved norm**

**Maintain Probabilistic Interpretation**

# Entangled Systems

$$\Psi(x, y; t) = \sum_{i,j} c_{ij}(t) \alpha_i(x) \beta_j(y)$$

**How do multi-particle systems evolve?**

$$\mathcal{H} \supset y\Phi\bar{\Psi}\Psi = (y\phi + \epsilon\langle\chi|\phi|\chi\rangle) \bar{\Psi}\Psi$$

$$\langle\chi|\phi|\chi\rangle = \int d^3x_1 d^3y_1 d\tau |\Psi(x_1, y_1; \tau)|^2 (G_R(t, x; \tau, x_1) + G_R(t, y; \tau, x_1) + G_R(t, x; \tau, y_1) + G_R(t, y; \tau, y_1))$$

**Additive form demanded by Polchinski - Natural in Field theory!**

# Gauge Theories and Gravitation

Linear QFT Lagrangian, Shift bosonic field by expectation value

To Path Integral, add:

$$e^{iS_0 + i \int d^4x \left( e^{\frac{(A_\mu + \epsilon_\gamma \langle \chi | A_\mu | \chi \rangle)}{1 + \epsilon_\gamma}} J^\mu + \epsilon_{\tilde{\gamma}} \langle \chi | F_{\mu\nu} | \chi \rangle F^{\mu\nu} \right)}$$

Background Field

Gravitation

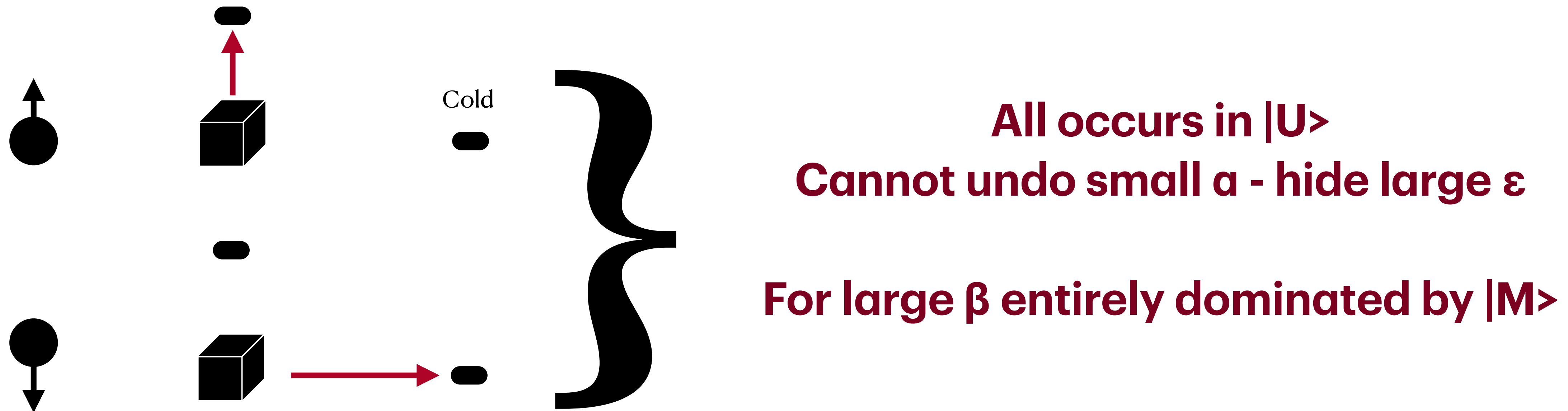
$$e^{iS_0 + i \int d^4x (\epsilon_G \langle \chi | g_{\mu\nu} | \chi \rangle \partial^\mu \phi \partial^\nu \phi)}$$

# Cosmological Sensitivity

# Non-Linearity and Cosmology

$$\mathcal{L} \supset e A_\mu \bar{\Psi} \gamma^\mu \Psi + \epsilon_\gamma e \langle \chi | A_\mu | \chi \rangle \bar{\Psi} \gamma^\mu \Psi$$

$$|\chi\rangle = \alpha|U\rangle + \beta|M\rangle \implies \langle \chi | A_\mu | \chi \rangle = |\alpha|^2 \langle U | A_\mu | U \rangle + |\beta|^2 \langle M | A_\mu | M \rangle$$



**Local Exploitability completely determined by unchangeable initial conditions**  
**Stark Difference from Linear Quantum Mechanics**

# Non-Linearity and Cosmology

$$\mathcal{L} \supset e A_\mu \bar{\Psi} \gamma^\mu \Psi + \epsilon_\gamma e \langle \chi | A_\mu | \chi \rangle \bar{\Psi} \gamma^\mu \Psi$$

$$|\chi\rangle = \alpha|U\rangle + \beta|M\rangle \implies \langle \chi | A_\mu | \chi \rangle = |\alpha|^2 \langle U | A_\mu | U \rangle + |\beta|^2 \langle M | A_\mu | M \rangle$$

**Could we boost by “Projection Operation”?**

$$\frac{\langle \chi | A_\mu | \chi \rangle}{\langle U | O | U \rangle}$$

**Environment is not local  
No local projection operator!**

**In linear quantum mechanics, we use this non-local projection operator - but degeneracy with coupling implies quantum phenomena understood without needing it**

**Likely a generic feature of local non-linear quantum mechanics**



# Inflationary Universe

**Are there effects that persist in the small  $a$  limit?**

$\langle \chi | \phi | \chi \rangle$  is homogeneous and isotropic BUT evolves in time

**Homogeneous fields will be non zero - but field needs to have a non-zero VEV**

**The Metric has a non-zero vev across the superposition!**

$$\tilde{g}_{\mu\nu} = \frac{g_{\mu\nu} + \epsilon_G \frac{\langle \chi | g_{\mu\nu} | \chi \rangle}{\langle \chi | \chi \rangle}}{1 + \epsilon_G}$$

# Inflationary Metric Interference

$$\tilde{g}_{\mu\nu} = \frac{g_{\mu\nu} + \epsilon_G \frac{\langle \chi | g_{\mu\nu} | \chi \rangle}{\langle \chi | \chi \rangle}}{1 + \epsilon_G}$$

**Inflation: Expectation value is the average homogeneous FRW metric**

$$g_s = - \left(1 - \frac{r_s}{r}\right) dt^2 + \frac{dr^2}{\left(1 - \frac{r_s}{r}\right)} + r^2 d\Omega^2 \qquad \langle g \rangle = -dt^2 + dr^2 + r^2 d\Omega^2$$

**Renormalize and Expand**

$$g_{\text{eff}} \simeq \left[ - \left(1 - \frac{R_s}{r}\right) dt^2 + \left(1 + \frac{R_s}{r} + \left(\frac{R_s}{r}\right)^2 (1 + \epsilon_G)\right) dr^2 \right] + r^2 d\Omega^2$$

**Long Distance Modification of Gravity!**

**Corrects Second Order GR Term!**

**Strong Field tests of GR!**

# Black Hole Horizon

$$\tilde{g}_{\mu\nu} = \frac{g_{\mu\nu} + \epsilon_G \frac{\langle \chi | g_{\mu\nu} | \chi \rangle}{\langle \chi | \chi \rangle}}{1 + \epsilon_G}$$

**Inflation: Expectation value is the average homogeneous FRW metric**

$$g_s = - \left(1 - \frac{r_s}{r}\right) dt^2 + \frac{dr^2}{\left(1 - \frac{r_s}{r}\right)} + r^2 d\Omega^2 \qquad \langle g \rangle = -dt^2 + dr^2 + r^2 d\Omega^2$$

$$g_{\text{eff}} = \frac{1}{1 + \epsilon_G} \left[ - \left(1 - \frac{r_s}{r} + \epsilon_G\right) dt^2 + \left( \frac{1}{1 - \frac{r_s}{r}} + \epsilon_G \right) dr^2 \right] + r^2 d\Omega^2$$

$g_{tt} \rightarrow 0, g_{rr} \rightarrow \infty$  at different values of  $r$ !

**Creates a firewall!**

# Conclusions

# Conclusions

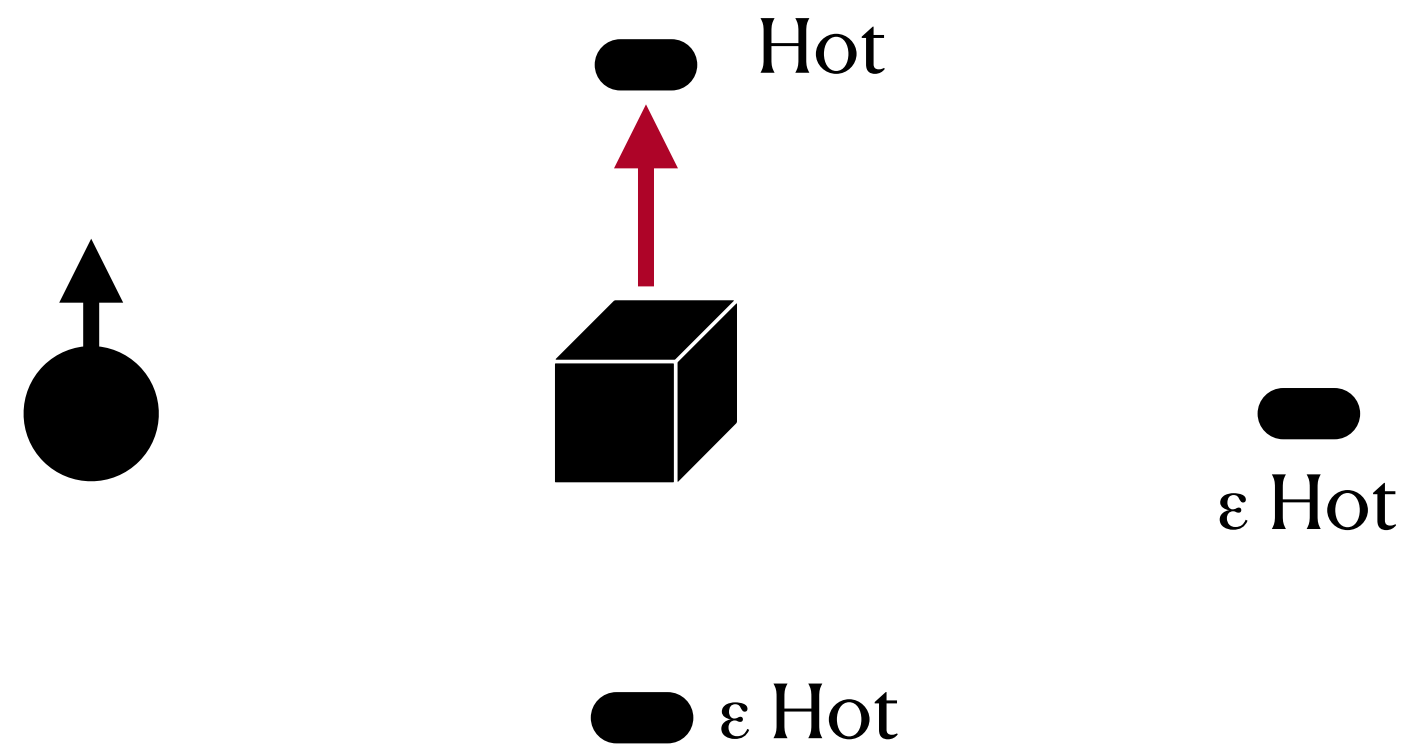
- 1. Quantum Field Theory can be generalized to include non-linear, state dependent time evolution**
- 2. Conventional tests of quantum mechanics in atomic and nuclear systems do NOT probe causal non-linear quantum mechanics**
- 3. Straightforward set of experimental tests possible to probe non-linear quantum mechanics**
- 4. Cosmological Sensitivity requires many more experimental probes**

**Backup**

# Delicate Non-Linearity

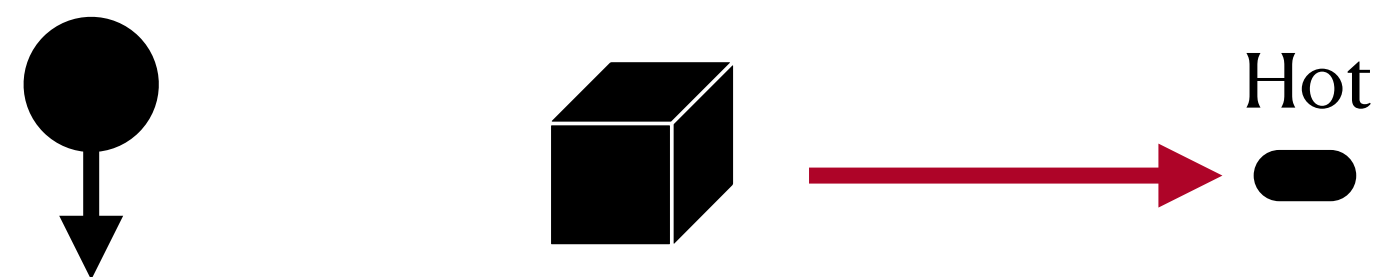
Suppose  $|X\rangle = |U\rangle$

O performs Laser experiment on July 6 - discovers non-linear quantum mechanics!



$$|\chi\rangle = \frac{1}{\sqrt{2}} (|U\rangle|O_U\rangle + |D\rangle|O_D\rangle)$$

O wants to repeat experiment



Suppose  $|O_U\rangle$  decides to run experiment at 9 AM on July 10  
But  $|O_D\rangle$  runs experiment on 9 AM on July 20

State on 9 AM on July 10

$$|\chi\rangle = \frac{1}{\sqrt{2}} \left( |U\rangle|O_U\rangle \frac{(|U\rangle|T\rangle + |D\rangle|R\rangle)}{\sqrt{2}} + |D\rangle|O_D\rangle \right)$$

# Delicate Non-Linearity

**State on 9 AM on July 10**

**Compare with State on July 6**

$$|\chi\rangle = \frac{1}{\sqrt{2}} \left( |U\rangle |O_U\rangle \frac{(|U\rangle |T\rangle + |D\rangle |R\rangle)}{\sqrt{2}} + |D\rangle |O_D\rangle \right)$$

$$|\chi\rangle = \frac{1}{\sqrt{2}} (|U\rangle |T\rangle + |D\rangle |R\rangle)$$

$$\langle U | \langle O_U | \langle U | \langle T | \langle E_T | e A_\mu (x_R) \bar{\Psi} (x_R) \gamma^\mu \Psi (x_R) | \chi \rangle = \frac{1}{2} \langle U | \langle T | \langle E_T | e A_\mu (x_R) \bar{\Psi} (x_R) \gamma^\mu \Psi (x_R) | \chi \rangle$$

**Effect is 1/2 of prior effect!**

**But, full effect if  $O_U$  and  $O_D$  perform experiment at same time!**

**Quantum Pollution: Without adequate care, superpositions may diverge wildly, preventing exploitability. Not automatic - but need careful protocols!**

**Particles have been scattering for 13 billion years. Cosmological dilution?**



# Cosmological Relaxation of Non-Linear QM?

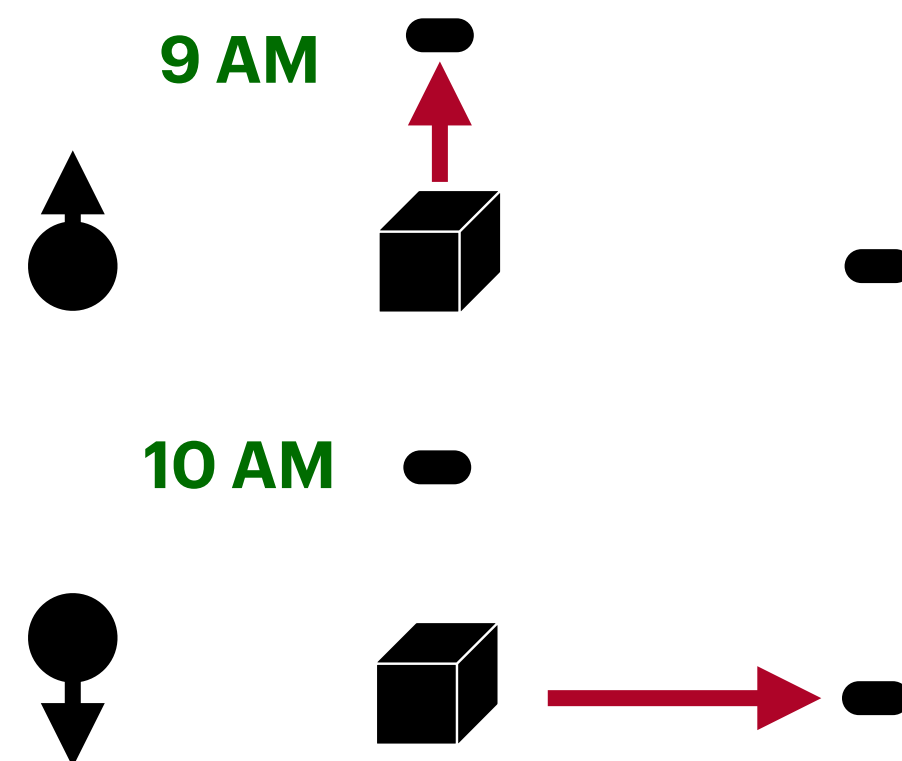
$$\mathcal{L} \supset e A_\mu \bar{\Psi} \gamma^\mu \Psi + \epsilon_\gamma e \langle \chi | A_\mu | \chi \rangle \bar{\Psi} \gamma^\mu \Psi$$

**All we need is the expectation value. Non-Linear effects are resistant to decoherence.**

**For e.g. when we repeat the experiment, it is ok for  $O_U$  and  $O_D$  to be two different individuals - all we care is that the fields are turned on at the same space-time points**

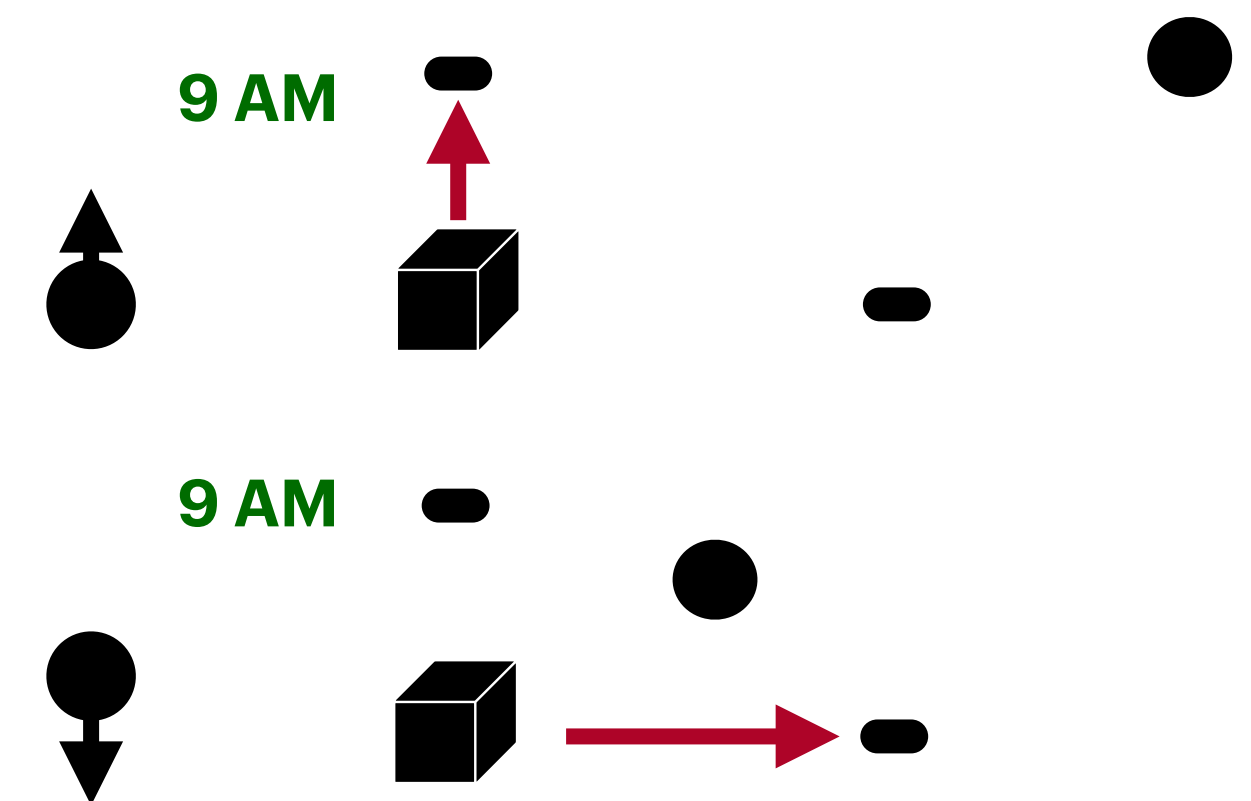
**Relevant**

**Superpositions where expectation values of fields are very different**



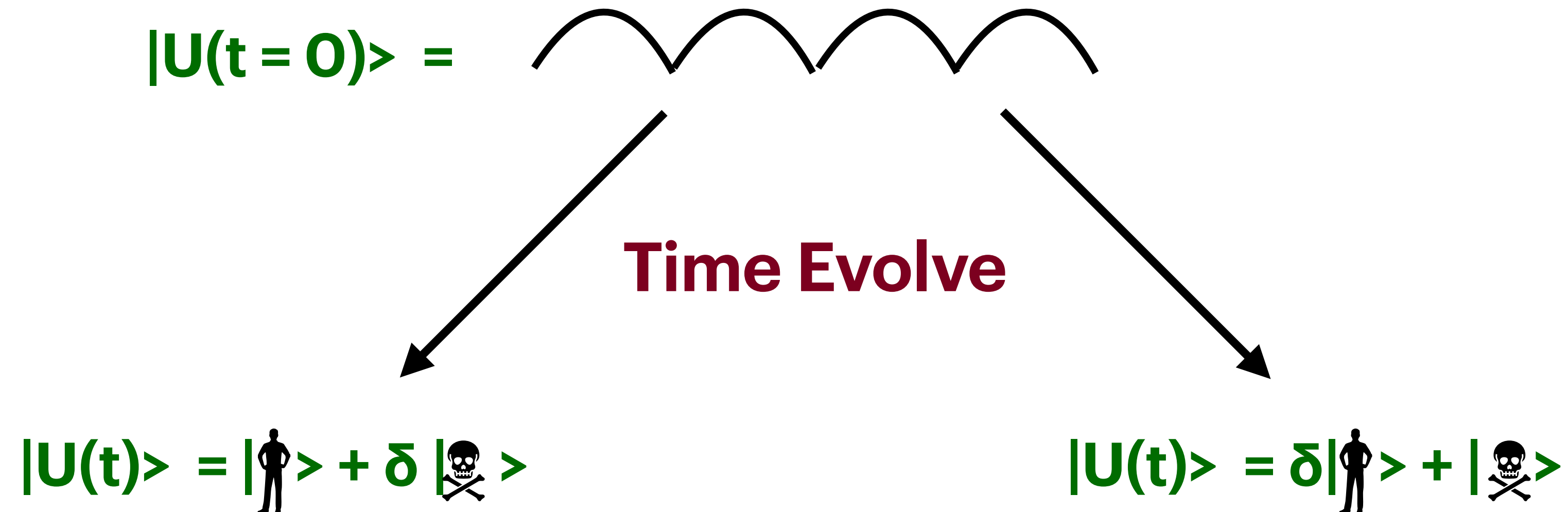
**Irrelevant**

**Scattering where expectation values are not significantly changed**



# Classical Universe?

Suppose  $|X\rangle = |U\rangle$



Can quantum events (scattering, decay etc.) lead to wildly different classical outcomes?

Clearly Possible - e.g. Human choosing to act differently based on quantum event

But, fundamentally - this is because humans can be quantum amplifiers

Are there natural quantum amplifiers, for e.g. in chaotic systems?

# Classical Universe?

$$|U(t)\rangle = |\text{person}\rangle + \delta |\text{skull}\rangle$$

Or

$$|U(t)\rangle = \delta |\text{person}\rangle + |\text{skull}\rangle$$

**Are there natural quantum amplifiers, for e.g. in chaotic systems?**

**Key Point: Changing classical evolution of a system requires coherent motion of N atoms**

**Probability that N atoms coherently move in some way:  $p^N$**

**With  $p \sim O(1)$  scattering probability**

**Even with  $N \sim 100$ , these are very small probabilities**

**For typical chaotic examples, e.g. Butterfly effect,  $N \gg 100$**

$$|U(t)\rangle = |\text{person}\rangle + \delta |\text{skull}\rangle$$

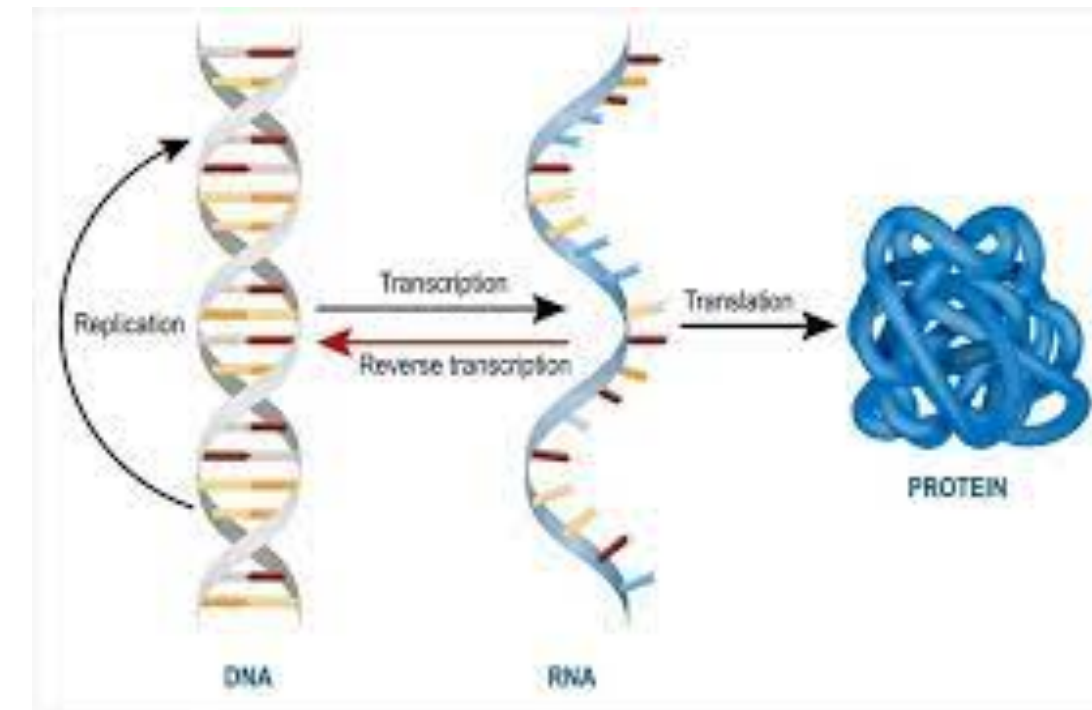
**Reasonable!**

**Quantum Amplifiers are Hard!**

# Evolutionary Dilution?

Humans can be quantum amplifiers

Is  $N \sim O(\text{few})$  for biology?  
Maybe for RNA/DNA?



$$|U(t=0)\rangle = \text{[wavy line symbol]}$$
$$|U(t)\rangle = |\text{[solar system icon]}\rangle (|\text{[human icon]}\rangle + |\text{[goat icon]}\rangle + \dots)$$

Dilutes Laboratory Effects!

# Inflation

**Quantum Amplifiers are Hard!**

**Except in Cosmic Inflation!**

**Inflation rapidly places quantum state in a homogenous and isotropic state (Bunch - Davies Vacuum )**

$\langle \chi | \phi | \chi \rangle$  **is homogeneous and isotropic BUT evolves in time**



**But, our universe is clearly inhomogeneous**  
**How could homogeneous state become inhomogeneous?**

**Answer: Massive Superposition of Statistically Similar Universes!**

$$|\chi\rangle = \sum_i c_i |U_i\rangle, \quad c_i \sim e^{-N}$$

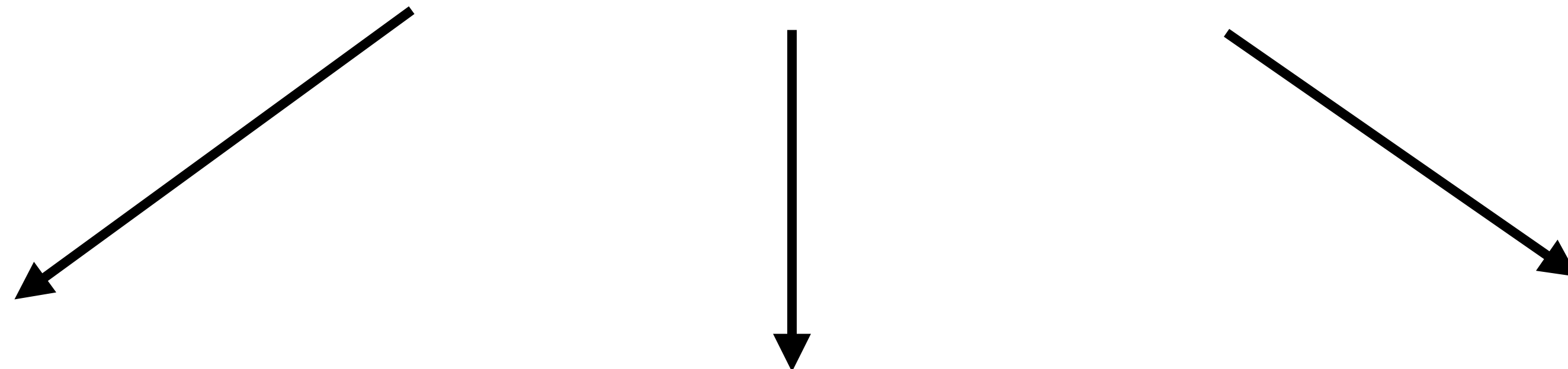
**Most of the Universe,**  
**The space-time point where the Earth is located is in intergalactic space!**

# Probes of Non-Linear Quantum Mechanics

Non-Linear effects deeply tied to unchangeable initial and evolution of total quantum state

Probe not just non-linear quantum mechanics but full evolutionary history of quantum state!

## Three Scenarios



### Classical Universe

$$|U(t)\rangle = |\text{person}\rangle + \delta |\text{skull and crossbones}\rangle$$

E.g. warm inflation,  
bouncing cosmologies

### Quantum Biology

$$|U(t)\rangle = |\text{universe}\rangle (|\text{person}\rangle + |\text{goat}\rangle + \dots)$$

### Canonical Inflation

$$|\chi\rangle = \sum_i c_i |U_i\rangle, \quad c_i \sim e^{-N}$$

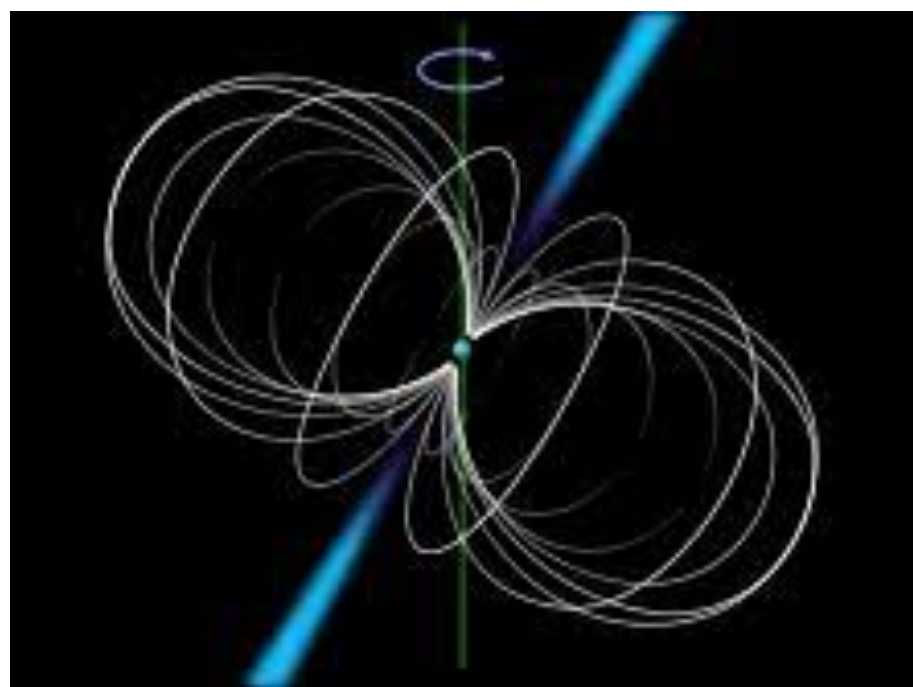


# Quantum Biology

$$|U(t)\rangle = | \text{🌌} \rangle (| \text{🧑} \rangle + | \text{🐐} \rangle + \dots)$$

The large scale structure of the universe is the same

Local structures (e.g. buildings) are vastly different



Shield

Look for **coherent** astrophysical or geological source (e.g. radio source/ magnetic field) inside shield!

$$\epsilon_\gamma \lesssim 10^{-5} \quad \text{Magnetic Shields Work!}$$

For gravity, perhaps look for gravitational effects of waves on man-made islands?

# Inflationary Universe

$$|\chi\rangle = \sum_i c_i |U_i\rangle, \quad c_i \sim e^{-N}$$

**Earth is in the middle of nowhere in the vast majority of the universes**

**Look for cosmic rays inside shields!**

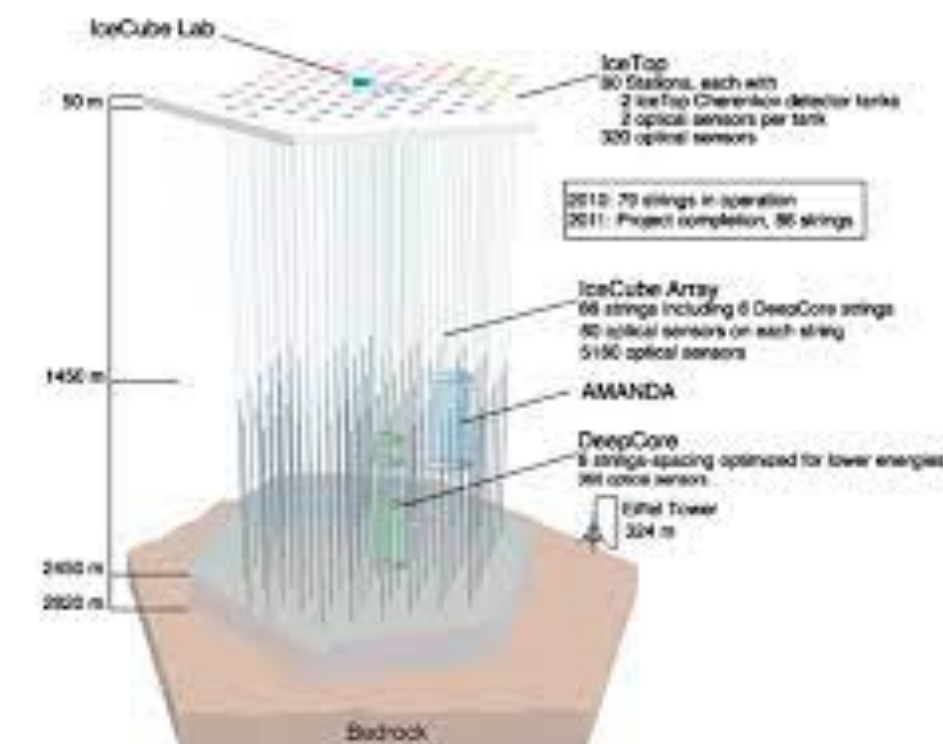


**Shield**

**e.g. Coherent Radio Sources**

$$\text{Current: } \propto \frac{1}{e^{\frac{N}{2}}}$$

**Maybe visible if N is not too big**



**e.g. protons in IceCube**